# **Vibrational Spectroscopies**

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# Vibrational Spectroscopies

"..everything that living things do can be understood in terms of the jigglings and wigglings of atoms.." R. P. Feymann

# **Vibrational Spectroscopies**

All vibrational spectroscopies are founded in interaction of functional groups in molecules with electro-magnetic radiation

What kinds of vibrational spectroscopies are outthere?

- Infrared spectroscopy
- Raman scattering
- High-resolution electron energy loss spectroscopy
- Non-linear vibrational spectroscopies (sum frequency generation, second harmonic generation)
- Scanning tunneling vibrational spectroscopy (performed by analyzing inelastic tunneling spectra taken with scanning electron microscope on a single molecule)



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#### Infrared spectroscopy

Infrared spectroscopy exploits the fact that molecules have specific frequencies at which they rotate or vibrate corresponding to discrete <u>energy levels</u> (vibrational modes). These <u>resonant frequencies</u> are determined by the shape of the molecular <u>potential</u> <u>energy surfaces</u>, the masses of the atoms and, by the associated <u>vibronic coupling</u>. In order for a vibrational mode in a molecule to be IR active, it must be associated with changes in the **permanent dipole**.

#### Infrared spectroscopy

The infrared portion of the electromagnetic spectrum is divided into three regions; the near-, mid- and far- infrared, named for their relation to the visible spectrum. The farinfrared, approximately 400–10 cm<sup>-1</sup> (1000–30 µm), lying adjacent to the <u>microwave</u> region, has low energy and may be used for <u>rotational spectroscopy</u>. The mid-infrared, approximately 4000–400 cm<sup>-1</sup> (30–2.5  $\mu$ m) may be used to study the fundamental vibrations and associated rotationalvibrational structure. The higher energy near-IR, approximately 14000–4000 cm<sup>-1</sup> (2.5–0.8 µm) can excite overtone or harmonic vibrations. The names and classifications of these subregions are merely conventions. They are neither strict divisions nor based on exact molecular or electromagnetic properties.

#### Infrared Spectroscopy



### Infrared Spectroscopy (Equipment)



#### Fourier transform infrared (FTIR) spectroscopy is a

measurement technique for collecting infrared spectra. Instead of recording the amount of energy absorbed when the frequency of the infra-red light is varied (monochromator), the IR light is guided through an <u>interferometer</u>. After passing through the sample, the measured signal is the interferogram. Performing a <u>Fourier transform</u> on this signal data results in a spectrum identical to that from conventional (dispersive) infrared spectroscopy.



(b)





Thin-film transmission spectra of the photocatalytic decomposition of ethanol: (a) before UV illumination, (b) after 3 h of UV illumination

Sho Kataoka, M. Isabel Tejedor-Tejedor, Juan M. Coronado, Marc A. Anderson Journal of Photochemistry and Photobiology A: Chemistry 163 (2004) 323–329

2) Attenuated total reflectance infrared spectroscopy (ATR-FTIR)





Using a pipette to add a liquid sample to a ZnSe HATR trough plate.



Applying pressure to a solid sample on the Universal diamond ATR top-plate.

http://las.perkinelmer.com/content/TechnicalInfo/TCH\_FTIRATR.pdf

2) Attenuated total reflectance infrared spectroscopy (ATR-FTIR)



Fiber-optic MolecularEye ATR diamond probe from



http://www.chemshow.cn/UploadFile/datum/1000/bosheng\_2009114141928791780.pdf

2) Attenuated total reflectance infrared spectroscopy (ATR-FTIR)

Multiple-internal reflectance Fourier-transform infrared spectroscopy (MIR-FTIR)



Bulanin, K.; Shah, A. and Teplyakov, A. V. Infrared spectroscopy studies of iodoethane on Si(100)-2x1: Adsorption and thermal decomposition leading to adsorbate ordering. *J. Chem. Phys.* **2001**, *115*(*15*), 7187-7195.



Wavenumber, cm<sup>-1</sup>

#### 3) Diffuse reflectance



Diffuse Reflectance Infrared Fourier Transform Spectroscopy (DRIFTS) is a technique that collects and analyzes scattered IR energy. It is used for measurement of fine particles and powders, as well as rough surface (e.g., the interaction of a surfactant with the inner particle, the adsorption of molecules on the particle surface). Sampling is fast and easy because little or no sample preparation is required.

• **Specular reflection**: When an infrared beam is focused on the surface of a particulate sample, it can interact with the sample two ways. First it may simply reflect off the sample surface in the same way visible light reflect off a mirror. This phenomenon, called "specular reflection", is a function of the refractive index of the sample.

• Specular reflection is suitable for smooth surface that can strongly reflect IR beam.

• **Diffuse reflection** collects infrared radiation from the IR beam that passes through a particle can either reflect off the next particle or be transmitted through the next particle. This transmission-reflectance event can occur many times in the sample, which increases the pathlength.

- Diffuse reflection collects scattered IR energy,
- Diffuse reflection is used for rough sample and particle

4) FTIR Imaging

#### • FTIR microscope

- Provide FT-IR spectra for samples down to about 10 µm in size
- Two primary modes of FT-IR microscopy: reflectance and transmission
- More advanced mode: ATR





#### 4) FTIR Imaging



**Raman spectroscopy** (named after <u>C. V. Raman</u>, pronounced (ro:mən) is a <u>spectroscopic</u> technique used to study vibrational, rotational, and other low-frequency modes in a system. It relies on <u>inelastic scattering</u>, or <u>Raman scattering</u>, of <u>monochromatic</u> light, usually from a <u>laser</u> in the <u>visible</u>, <u>near infrared</u>, or <u>near ultraviolet</u> range. The laser light interacts with <u>phonons</u> or other excitations in the system, resulting in the energy of the laser photons being shifted up or down. The shift in energy gives information about the phonon modes in the system. <u>Infrared spectroscopy</u> yields similar, but complementary, information.

The Raman effect occurs when light impinges upon a <u>molecule</u> and interacts with the <u>electron cloud</u> and the <u>bonds</u> of that molecule. For the spontaneous Raman effect, a <u>photon</u> excites the molecule from the <u>ground state</u> to a virtual <u>energy state</u>. When the molecule emits a photon and returns to the ground state, it returns to a different <u>rotational</u> or <u>vibrational state</u>. The difference in energy between the original state and this new state leads to a shift in the emitted photon's frequency away from the excitation frequency.

If the final state of the molecule is more energetic than the initial state, then the emitted photon will be shifted to a lower frequency in order for the total energy of the system to remain balanced. This shift in frequency is designated as a <u>Stokes shift</u>. If the final state is less energetic than the initial state, then the emitted photon will be shifted to a higher frequency, and this is designated as an anti-Stokes shift. Raman scattering is an example of inelastic scattering because of the energy transfer between the photons and the molecules during their interaction.

A change in the molecular polarization potential — or amount of deformation of the electron cloud — with respect to the vibrational coordinate is required for a molecule to exhibit a Raman effect. The amount of the polarizability change will determine the Raman scattering intensity. The pattern of shifted frequencies is determined by the rotational and vibrational states of the sample.



The basic set-up of a Raman spectrometer is shown on the right. Note that the detector is orthogonal to the direction of the incident radiation, so as to observe only the scattered light. The source needs to provide intense monochromatic radiation, it is usually a laser.

1) Raman spectroscopy collects scattered radiation.

2) The criteria for a molecule to be Raman active are also different from other types of spectroscopy. It does NOT require a permanent dipole moment.



#### Selection rule for Raman activity -- polarizability

For IR spectroscopy, it is necessary for the molecule to have a permanent electric dipole. This is not the case for Raman spectroscopy, rather it is the polarizability ( $\alpha$ ) of the molecule which is important.

The oscillating electric field of a photon causes charged particles (electrons and, to a lesser extent, nuclei) in the molecule to oscillate. This leads to an induced electric dipole moment,  $\mu_{ind}$ , where

#### $\mu_{ind} = \alpha E$

This induced dipole moment then emits a photon, leading to either Raman or Raleigh scattering.

The energy of this interaction is also dependent on the polarizability:

Energy of interaction =  $-1/2\alpha E^2$ 

The energies of Raman transitions are relatively weak. To counter this, a higher intensity of the exciting radiation is used.

For Raman scattering to occur, the polarizability of the molecule must vary with its orientation.

IR	Raman
I. Vibrational Modes	Vibrational Modes
II. Change in dipole	II. Change in polarizability
$\overbrace{\delta}^{\rightarrow} \underbrace{-2\delta^{+}}_{2\delta^{+}} \underbrace{-\delta}_{\delta}^{\rightarrow}$	extend
III. Excitation of a molecule to the excited vibrational state	III. Momentary distortion of the electron distribution

IR	Raman
IV. Absorption: (radiation at a certain frequency is absorbed due to the resonance of molecular vibration)	IV. (radiation at a certain frequency is scattered by the molecule with <i>shifts</i> in the wavelength of the incident beam)
V. Observed peak is due to molecular vibration	V. observed frequency shifts are related to vibrational changes in the molecule

The molecule is originally at the  $E_2$  energy state. The photon interacts with the molecule, exciting it with an energy  $hv_i$ . However, there is no stationary state of the molecule corresponding to this energy, and so the molecule relaxes down to the energy levels  $E_1$  or  $E_3$ . This process emits a photon in two ways:

Anti-Stokes transition:

If the molecule relaxes to energy state  $E_1$ , it will have lost energy, and so the photon emitted will have energy  $hv_{r1}$ , where  $hv_{r1} > hv_i$ 

Stokes transitions:

If the molecule relaxes to energy state  $E_3$ , it will have gained energy, and so

the photon emitted will have energy  $hv_{r2}$ , where  $hv_{r2} < hv_i$ 

For rotational spectroscopy, both Stokes and anti-Stokes transitions are observed.

For vibrational spectroscopy, Stokes transitions are far more common, and so anti-Stokes transitions can effectively be ignored.







## Flow Field Plate - Graphite $\lambda_{ex} = 514.5 \text{ nm}$ Intensity (counts) $\left(\frac{I_g}{I_s}\right) = 3.98$ 500 1000 1500 2000 2500 0 Raman Shift, $\Delta \sigma$ (cm<sup>-1</sup>)



Nanocrystalline graphite has graphitic (g) and disorder (d) peaks. The characteristic dimension of graphitic domains is given by:

$$\square \hat{A}) = 44 \left(\frac{I_g}{I_d}\right)$$
$$\square = 175 \hat{A}$$

From early literature on graphitic materials Tuinstra and Koenig, J. Chem Phys. 53, 1126 (1970).





Confocal Raman microscopy for three-dimensional imaging







#### Electron Energy Loss Spectroscopy (EELS)

In electron energy loss spectroscopy (EELS) a material is exposed to a beam of electrons with a known, narrow range of kinetic energies. Some of the electrons will undergo inelastic scattering, which means that they lose energy and have their paths slightly and randomly deflected. The amount of energy loss can be measured via an electron spectrometer and interpreted in terms of what caused the energy loss. Inelastic interactions include phonon excitations, inter and intra band transitions, plasmon excitations, inner shell ionisations, and Cherenkov radiation. The inner shell ionizations are particularly useful for detecting the elemental components of a material. For example, one might find that a larger-than expected number of electrons comes through the material with 285 eV (electron volts, a unit of energy) less energy than they had when they entered the material. It so happens that this is about the amount of energy needed to remove an inner-shell electron from a carbon atom. This can be taken as evidence that there's a significant amount of carbon in the part of the material that's being hit by the electron beam. With some care, and looking at a wide range of energy losses, one can determine the types of atoms, and the numbers of atoms of each type, being struck by the beam. The scattering angle (that is, the amount that the electron's path is deflected) can also be measured, giving information about the dispersion relation of whatever material excitation caused the inelastic scattering.



Based on inelastic scattering of monoenergetic beam of low energy electrons ( $E_i = 1-10 \text{ eV}$ ) from surface

- kinetic energy analysis of inelastically-scattered electrons gives information about energy transfer to surface vibrational modes

-  $E_s = E_i - hv$  --- energy loss

 $H_2$  adsorbed on TM expect  $hv = 4560 \text{ cm}^{-1}$  but for H on TM expect  $hv = 800-1000 \text{ cm}^{-1}$ 

- minimum resolution  $\Delta E_i$  (FWHM of incident beam) of EELS > RAIRS - best <8 cm<sup>-1</sup>, typically 20-30 cm<sup>-1</sup> (8 cm<sup>-1</sup> @ 1 meV)



Electron Scattering Mechanisms in EELS Three basic scattering modes for electron, each with own selection rules:

- (1) Dipole scattering
- (2) Impact scattering
- (3) Negative ion resonance scattering
- gives EELS added versatility over RAIRS
- all three modes operative (one usually dominates under certain experimental conditions)

#### (1) Dipole scattering

Electron is scattered at long range (100's Å) by oscillating Coulombic field of surface

Electron momentum  $\perp$  to surface is conserved

Energy loss results in small angle deflections around specular direction - dipole scattering lobe

- Elastically-scattered electrons appear at  $\Theta_i = \Theta_s$  (specular)
- Inelastically-scattered electrons appear at  $\Theta_i \approx \Theta_s$  (sharply peaked near specular)





small momentum loss  $q_{\parallel} = hv$ 

Only modes with component of vibration perpendicular to surface can be excited (not parallel modes) - *surface selection rule* 

same selection rules as RAIRS



Large cross-section for small  $E_i\mbox{'s}\ (\mbox{<}20\ eV)$  - dominant mechanism near specular

Well-developed/straightforward theory

#### (2) Impact scattering:

Electron is "kinematically" scattered at short range (few Å) by atomic core potentials of surface - modulated by nuclear motion (vibration)

Electron scattered over wide range of angles (in and out of plane of incidence)

Probability of scattering depends on surface dipole amplitude/direction (symmetry of mode) and electron energy (momentum)

Neither electron momentum  $\parallel$  nor  $\perp$  conserved

- can see vibrational modes with dynamic dipoles both  $\parallel$  and  $\perp$  to surface plane (only off-specular)

Much lower cross-section than dipole scattering (weak loss features)

Poorly-developed/complex theory

(3) Negative Ion Resonance

Incident electron becomes transiently trapped in empty, high lying MO of adsorbate - Rydberg state

Vibrational feature intensity strongly depends on incident energy - resonances

Very small cross-section off-resonance (rare)

Regardless of mechanism, most electrons scattered *elastically* - loss spectrum usually very weak

## The Electron Source

The electron beam comes from a filament, made of various types of materials. The most common is the Tungsten hairpin gun. This filament is a loop of tungsten which functions as the cathode. A voltage is applied to the loop, causing it to heat up. The anode, which is positive with respect to the filament, forms powerful attractive forces for electrons. This causes electrons to accelerate toward the anode. Some accelerate right by the anode and on down the column, to the sample. Other examples of filaments are Lanthanum Hexaboride filaments and field emission guns.



#### Instrumentation for HREELS

Monochromation of electrons emitted by thermionic emission achieved by passing through spherical-sector electron energy analyzer

typical radius 5 cm, typical pass energy <500 meV (in XPS, radius ~20-30 cm, pass energy 50-100 eV)</li>

Electrons retarded before entry into sector, then accelerated to sample

Second spherical-sector used to analyze electron kinetic energy

Entire instrument magnetically shielded - no magnetic materials for construction

In very high resolution instruments, two sectors for monochromator and two for analyzer



#### Instrumentation for HREELS



Space-charge effects always tend to produce divergent beam

Optimization of lens and sector potentials very difficult - computer controllec

In best HREELS, analyzer rotatable for off-specular (impact scattering) analysis

E<sub>i</sub> variable 1-50 eV or more - can study vibrational and electronic loss processes (plasmons)

- observe electronic excitation from 0->50 eV
- observe vibrational modes from 0->4000 cm<sup>-1</sup>



Fig. 9.11 Set of imaginary adsorption states showing the expected loss peaks associated with each structure.

#### Example Application of HREELS: CO adsorption on W(100)



At low CO exposures

- 70 meV (565 cm<sup>-1</sup>) W-C (hollow site)
- 78 meV (630 cm<sup>-1</sup>) W-O (hollow site)

#### At high pCO

- 45 meV (363 cm<sup>-1</sup>) W-CO (on-top)
- $258 \text{ meV} (2081 \text{ cm}^{-1}) \text{ -C=O} (CO(g) = 2140 \text{ cm}^{-1})$

Fig. 9.13 Electron energy loss spectra of CO on W{100} at 300 K. The spectra are recorded with a primary electron beam energy of 5 eV incident at an angle of 75°;  $\Delta \phi$  is the measured increase in work function (Froitzheim *et al.*, 1976).



Figure 2.25. Vibrational spectra from chemisorbed ethylene on Rh(111) at different temperatures obtained by HREELS. Note the sequential dehydrogenation process.

At 77 K, ethylene non-dissociatively chemisorbs on Rh(111) with C=C approximately parallel to surface in  $\pi$ -bonded configuration



**Figure 7.13**. RAIRS spectrum of a monlayer of CO adsorbed on a Cu (111) surface at 95 K in comparison with analogous data obtained using electron energy loss spectroscopy, after Chesters *et al.* [10]

## Summary:

"Fingerprint technique" - can identify functional groups through vibrational frequency analysis

Use of selection rules/group theory allows determination of bond orientation

Difficult to quantify species (RAIRS easier then HREELS)

#### RAIRS:

Excellent energy resolution (<2 cm<sup>-1</sup>) - useful for separating multiple peaks, phase transitions, lateral interactions, dynamics of coupling

Straightforward instrumentation

Not restricted to surfaces in vacuum - can be used in "real world" conditions from several atm to UHV

Theory well-developed

Rapid - 30 s-10 minutes - time-resolved spectra possible

BUT

Less sensitive than HREELS

Cannot observe modes <800 cm<sup>-1</sup>

Only dipole active modes - component perpendicular to surface

Need background (reference spectrum) subtraction

#### HREELS:

Multiple scattering mechanisms allow observation of modes parallel and perpendicular to surface

Observe modes 0->4000 cm<sup>-1</sup> - low frequency modes (molecule-surface) visible

Can perform vibrational and electronic loss spectroscopy

Can induce and probe electron-induced reactions with variable E<sub>i</sub>

No background subtraction

Most common technique for surface and bulk phonon measurement

BUT

Poor resolution (6-8 cm<sup>-1</sup> best)

Only works in UHV

Very expensive, complex, delicate instrumentation

Slow - 15-60 minutes per spectrum

Electron spectroscopy - surface charging on insulators

Difficult theory